

# Long-time behavior of solutions to some equations in fluid mechanics

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Hanoi, August 19, 2014

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# Some basic problems for PDEs

$$\begin{cases} \partial_t u = \mathcal{F}(t, u(t)), t > \tau, \\ u(\tau) = u_\tau. \end{cases}$$

- Well-posedness
- Regularity of solutions
- Long-time behavior of solutions: Stability theory + attractors theory
- Control theory:
  - Controllability
  - Optimal control
  - Stabilization

# Process

$$\begin{cases} \partial_t u = \mathcal{F}(t, u(t)), t > \tau, \\ u(\tau) = u_\tau \in X. \end{cases} \quad (1.1)$$

Assume that for each  $u_\tau \in X$  given, problem (1.1) has a unique global solution  $u(t; \tau, u_\tau)$ . Putting  $U(t, \tau)u_\tau := u(t; \tau, u_\tau)$ , we get a process  $U(t, \tau)$  on  $X$ , which is called the process associated to problem (1.1).

A process on  $X$  is a family of two-parameter mappings  $\{U(t, \tau)\}$  in  $X$  having the following properties:

$$\begin{aligned} U(\tau, \tau) &= Id \text{ for all } \tau \in \mathbb{R}, \\ U(t, r)U(r, \tau) &= U(t, \tau) \text{ for all } t \geq r \geq \tau. \end{aligned}$$

# Theory of attractors

Study the long-time behavior of solutions: theory of attractors

- Autonomous PDEs: Solution semigroup  $S(t) : u_0 \mapsto u(t)$ .  
Since 1980s: Theory of global attractors.
- Non-autonomous PDEs: The associated process:

$$U(t, \tau) : u_\tau \mapsto U(t, \tau)u_\tau = u(t).$$

- Theory of uniform attractors: Chepyzhov-Vishik (1994);
- Theory of pullback attractors: Caraballo-Łukaszewicz-Real (2006).
- Advantages of pullback attractors:
  - allow to handle a larger class of time-dependent external forces;
  - (Usually) have a finite fractal dimension;
  - are also valid for random dynamical systems.

# Definition of pullback $\mathcal{D}$ -attractors

Suppose that  $\mathcal{B}(X)$  is the family of all nonempty bounded subsets of  $X$ , and  $\mathcal{D}$  is a non-empty class of parameterized sets

$$\hat{\mathcal{D}} = \{D(t) : t \in \mathbb{R}\} \subset \mathcal{B}(X).$$

A family  $\hat{A} = \{A(t) : t \in \mathbb{R}\} \subset \mathcal{B}(X)$  is said to be a pullback  $\mathcal{D}$ -attractor for the process  $U(t, \tau)$  if

- ①  $A(t)$  is compact for all  $t \in \mathbb{R}$ ;
- ②  $\hat{A}$  is invariant, i.e.,  $U(t, \tau)A(\tau) = A(t)$ , for all  $t \geq \tau$ ;
- ③  $\hat{A}$  is pullback  $\mathcal{D}$ -attracting, i.e.,

$$\lim_{\tau \rightarrow -\infty} \text{dist}(U(t, \tau)D(\tau), A(t)) = 0, \text{ for all } \hat{D} \in \mathcal{D}, \text{ and all } t \in \mathbb{R};$$

- ④ If  $\{C(t) : t \in \mathbb{R}\}$  is another family of closed attracting sets, then  $A(t) \subset C(t)$ , for all  $t \in \mathbb{R}$ .

## Some basic concepts

- The process  $U(t, \tau)$  is said to be pullback  $\mathcal{D}$ -asymptotically compact if for any  $t \in \mathbb{R}$ , any  $\hat{D} \in \mathcal{D}$ , any sequence  $\tau_n \rightarrow -\infty$ , and any sequence  $x_n \in D(\tau_n)$ , the sequence  $\{U(t, \tau_n)x_n\}$  is relatively compact in  $X$ .
- The family of bounded sets  $\hat{B} \in \mathcal{D}$  is called pullback  $\mathcal{D}$ -absorbing for the process  $U(t, \tau)$  if for any  $t \in \mathbb{R}$ , any  $\hat{D} \in \mathcal{D}$ , there exists  $\tau_0 = \tau_0(\hat{D}, t) \leq t$  such that

$$\bigcup_{\tau \leq \tau_0} U(t, \tau)D(\tau) \subset B(t).$$

# Existence of pullback $\mathcal{D}$ -attractors

Let  $\{U(t, \tau)\}$  be a continuous process on  $X$  such that

- ①  $\{U(t, \tau)\}$  is pullback  $\mathcal{D}$ -asymptotically compact;
- ② there exists a family of pullback  $\mathcal{D}$ -absorbing sets  $\hat{\mathcal{B}} = \{B(t) : t \in \mathbb{R}\} \in \mathcal{D}$ .

Then  $\{U(t, \tau)\}$  has a unique pullback  $\mathcal{D}$ -attractor  $\hat{\mathcal{A}} = \{A(t) : t \in \mathbb{R}\}$ , and

$$A(t) = \bigcap_{s \leq t} \overline{\bigcup_{\tau \leq s} U(t, \tau) B(\tau)}.$$



## Some references

- 1 T. Caraballo, G. Łukaszewicz and J. Real, Pullback attractors for asymptotically compact non-autonomous dynamical systems, *Nonlinear Anal.* 64 (2006), 484-498.
- 2 A.N. Carvalho, J.A. Langa and J.C. Robinson, *Attractors for Infinite-Dimensional Non-autonomous Dynamical Systems*, Applied Mathematical Sciences 182, Springer, Berlin, 2013.

# Navier-Stokes-Voigt equations

Let  $\Omega$  be a (bounded or unbounded) domain in  $\mathbb{R}^d$  ( $d = 2, 3$ ) with boundary  $\partial\Omega$ . Consider the following problem

$$\left\{ \begin{array}{l} \partial_t u - \nu \Delta u - \alpha^2 \Delta u_t + (u \cdot \nabla)u + \nabla p \\ \nabla \cdot u \\ u(x, t) \\ u(x, \tau) \end{array} \right. \begin{array}{l} = f, \quad x \in \Omega, t > \tau, \\ = 0, \quad x \in \Omega, t > \tau, \\ = 0, \quad x \in \partial\Omega, t > \tau, \\ = u_0(x), \quad x \in \Omega, \end{array} \quad (2.1)$$

where  $u = u(x, t) = (u_1, \dots, u_d)$  is the unknown velocity vector,  $p = p(x, t)$  is the unknown pressure,  $\nu > 0$  is the kinematic viscosity coefficient, and  $u_0$  is the initial velocity.

- Meaning: a model of linear viscoelastic fluids (Oskolkov, 1973) + a regularization of the Navier-Stokes equations (Titi et. al., 2006).
- Difficulties: the system is only weakly dissipative + the domain may be unbounded.

# Hypotheses

- (H1) The domain  $\Omega$  can be an arbitrary (bounded or unbounded) domain in  $\mathbb{R}^d$  without any regularity assumption on its boundary  $\partial\Omega$ , provided that the Poincaré inequality holds on  $\Omega$ : There exists  $\lambda_1 > 0$  such that

$$\int_{\Omega} \phi^2 dx \leq \frac{1}{\lambda_1} \int_{\Omega} |\nabla \phi|^2 dx, \quad \forall \phi \in H_0^1(\Omega).$$

- (H2)  $f \in L_{loc}^2(\mathbb{R}; V')$  such that

$$\int_{-\infty}^0 e^{\sigma s} \|f(s)\|_{V'}^2 ds < +\infty,$$

where  $\sigma = \frac{\lambda_1 \nu}{1 + \alpha^2 \lambda_1}$ .

# Function spaces

Let  $(L^2(\Omega))^d$  and  $(H_0^1(\Omega))^d$  be endowed with the inner products

$$(u, v) = \int_{\Omega} u \cdot v dx, \quad u, v \in (L^2(\Omega))^d,$$

$$((u, v)) = \int_{\Omega} \sum_{j=1}^d \nabla u_j \cdot \nabla v_j dx, \quad u, v \in (H_0^1(\Omega))^d,$$

and norms  $|u|^2 = (u, u)$ ,  $\|u\|^2 = ((u, u))$ .

Let

$$\mathcal{V} = \{u \in (C_0^\infty(\Omega))^d : \nabla \cdot u = 0\}.$$

Denote by  $H$  the closure of  $\mathcal{V}$  in  $(L^2(\Omega))^d$ , and by  $V$  the closure of  $\mathcal{V}$  in  $(H_0^1(\Omega))^d$ . It follows that  $V \subset H \equiv H' \subset V'$ , where the injections are dense and continuous. We will use  $\|\cdot\|_*$  for the norm in  $V'$ , and  $\langle \cdot, \cdot \rangle$  for the duality pairing between  $V$  and  $V'$ .

# Operators

We now define the trilinear form  $b$  by

$$b(u, v, w) = \sum_{i,j=1}^d \int_{\Omega} u_i \frac{\partial v_j}{\partial x_i} w_j dx,$$

whenever the integrals make sense. It is easy to check that if  $u, v, w \in V$ , then

$$b(u, v, w) = -b(u, w, v).$$

Hence

$$b(u, v, v) = 0, \quad \forall u, v \in V.$$

Set  $A : V \rightarrow V'$  by  $\langle Au, v \rangle = ((u, v))$ ,  $B : V \times V \rightarrow V'$  by  $\langle B(u, v), w \rangle = b(u, v, w)$ ,  $Bu = B(u, u)$ . Then  $Au = -P\Delta u$ , for all  $u \in D(A)$ , where  $P$  is the ortho-projector from  $(L^2(\Omega))^d$  onto  $H$ .

# Basic inequalities

Using Hölder's inequality and Ladyzhenskaya's inequality:

$$|u|_{L^4} \leq c_0 |u|^{d/4} |\nabla u|^{1-d/4}, \quad \forall u \in H_0^1(\Omega),$$

and the interpolation inequalities, one can prove the following

- If  $d = 2, 3$ , then

$$|b(u, v, w)| \leq C |u|^{1-d/2} \|u\|^{d/4} |v|^{1-d/4} \|v\|^{d/4} \|w\|, \quad \forall u, v, w \in V.$$

In particular,

$$|b(u, u, v)| \leq C |u|^{2-2/d} \|u\|^{d/2} \|v\|, \quad \forall u, v \in V.$$

- Let  $u \in L^2(\tau, T; V)$ . Then the function  $Bu$  defined by

$$(Bu(t), v) = b(u(t), u(t), v), \quad \forall u \in V, \text{ a.e. } t \in [\tau, T],$$

belongs to  $L^1(\tau, T; V')$ .

## Definition of weak solutions

Given  $u_0 \in V$ . A function  $u$  is called a **weak solution** to problem (2.1) on the interval  $(\tau, T)$  if

$$\begin{cases} u \in C([\tau, T]; V), \quad du/dt \in L^2(\tau, T; V), \\ \frac{d}{dt}u(t) + \nu Au(t) + \alpha^2 Au'(t) + B(u(t), u(t)) = f(t) \text{ in } V', \text{ for a.e. } t, \\ u(\tau) = u_0. \end{cases}$$

# Existence theorem

Suppose that  $u_0 \in V$  is given and assumptions (H1) – (H2) hold. Then, for any  $\tau \in \mathbb{R}$ ,  $T > \tau$  given, problem (2.1) has a unique weak solution  $u$  on  $(\tau, T)$ . Moreover, the map  $u_0 \mapsto u(t)$  is continuous on  $V$  for all  $t \in [\tau, T]$ .



# Sketch of the proof

Using the Galerkin approximation + the Compactness Lemma

- **Step 1:** Construct the approximate solutions  $u_m$

$$\begin{cases} \frac{d}{dt}u_m(t) + \nu Au_m(t) + \alpha^2 Au'_m(t) + P_m Bu_m(t) = P_m f(t) \text{ in } V', \\ u_m(\tau) = P_m u_0. \end{cases}$$

- **Step 2:** Establish some *a priori* estimates for  $u_m$

$$\begin{aligned} & |u^m(t)|^2 + \nu \int_{\tau}^t |\nabla u^m(s)|^2 ds + \alpha^2 |\nabla u^m(t)|^2 \\ & \leq \frac{1}{\nu} \|f\|_{L^2(\tau, T; V')}^2 + |u_0|^2 + \alpha^2 |\nabla u_0|^2. \end{aligned}$$

- **Step 3:** Passage to the limit
- **Step 4:** Prove the uniqueness and continuous dependence of the weak solutions on the initial data.

# The associated process

Thanks to the existence theorem, one can define a continuous process  $U(t, \tau)$  in  $V$  by

$$U(t, \tau)u_0 = u(t; \tau, u_0), \quad \tau \leq t, u_0 \in V,$$

where  $u(t) = u(t; \tau, u_0)$  is the unique weak solution of problem (2.1) with the initial datum  $u(\tau) = u_0$ .

**Lemma.** Let  $\{u_{0_n}\}$  be a sequence in  $V$  converging weakly in  $V$  to an element  $u_0 \in V$ . Then

$$U(t, \tau)u_{0_n} \rightharpoonup U(t, \tau)u_0 \quad \text{weakly in } V, \quad \text{for all } \tau \leq t,$$

$$U(t, \tau)u_{0_n} \rightharpoonup U(t, \tau)u_0 \quad \text{weakly in } L^2(\tau, T; V), \quad \text{for all } \tau < T.$$

# Existence of a pullback $\mathcal{D}$ -attractor

Let  $\mathcal{R}_\sigma$  be the set of all functions  $r : \mathbb{R} \rightarrow (0, +\infty)$  such that

$$\lim_{t \rightarrow -\infty} e^{\sigma t} r^2(t) = 0,$$

where  $\sigma = \frac{\lambda_1 \nu}{1 + \alpha^2 \lambda_1}$ , and  $\mathcal{D}_\sigma$  the class of all families

$\hat{\mathcal{D}} = \{D(t) : t \in \mathbb{R}\} \subset \mathcal{B}(H)$  such that  $D(t) \subset B(0, \hat{r}(t))$ , for some  $\hat{r}(t) \in \mathcal{R}_\sigma$ , where  $B(0, r)$  denotes the close ball in  $V$ , centered at zero with radius  $r$ .

**Theorem.** Suppose that conditions (H1) – (H2) hold. Then, there exists a unique pullback  $\mathcal{D}_\sigma$ -attractor  $\hat{A}_\alpha = \{A_\alpha(t) : t \in \mathbb{R}\}$  for the process  $\{U(t, \tau)\}$  associated to problem (2.1).

# Sketch of the proof

We will check the two conditions in the abstract theorem:

- 1  $U(t, \tau)$  has a family  $\hat{B}$  of pullback  $\mathcal{D}_\sigma$ -absorbing sets

$$\|u(t)\|^2 \leq \frac{1}{\alpha^2} e^{-\sigma(t-\tau)} [u_0]_2^2 + \frac{e^{-\sigma t}}{\nu \alpha^2} \int_\tau^t e^{\sigma r} \|f(r)\|_*^2 dr,$$

where  $[u]_2^2 := |u|^2 + \alpha^2 \|u\|^2$ . Denote

$R_\sigma^2(t) := \frac{2e^{-\sigma t}}{\nu} \int_{-\infty}^t e^{\sigma s} \|f(s)\|_*^2 ds$ , then  $\{B_\sigma(t) : t \in \mathbb{R}\}$  is a family of bounded pullback  $\mathcal{D}_\sigma$ -absorbing sets.

- 2  $U(t, \tau)$  is pullback  $\mathcal{D}_\sigma$ -asymptotically compact: Using the energy equation method introduced by J.M. Ball, *Disc. Cont. Dyn. Syst.* (2004).

# Estimates of the fractal dimension

We suppose furthermore that

$$f \in L^\infty(-\infty, T^*; V'), \quad \text{for some } T^* \in \mathbb{R}. \quad (2.2)$$

**Theorem.** Suppose that conditions (H1), (H2) and (2.2) hold. Then, the pullback  $\mathcal{D}_\sigma$ -attractor  $\hat{A}_\alpha = \{A_\alpha(t) : t \in \mathbb{R}\}$  of the process  $U(t, \tau)$  associated to problem (2.1) satisfies

$$d_F(A_\alpha(t)) \leq 2 + \frac{C(\lambda_1 + \alpha^2)^2 \|f\|_{L^\infty(-\infty, T^*; V')}^4}{\nu^6 \alpha^6 \sigma^2}, \quad \text{for all } t \in \mathbb{R}.$$

For more details, see Anh-Trang, *Proc. R. Soc. Edinb. Sect. A Math.* (2013).

# Regularity of the pullback attractor for Navier-Stokes-Voigt equations

We assume that

**(F)**  $f \in L^2_{loc}(\mathbb{R}; H)$ ,  $f' \in L^2_{loc}(\mathbb{R}; V')$  such that

$$\int_{-\infty}^0 e^{\frac{\sigma s}{3}} \|f(s)\|_{-1/2}^2 ds < +\infty;$$

$$\int_{-\infty}^0 e^{\sigma s} |f(s)|^2 ds < +\infty,$$

where  $\sigma = \frac{\lambda_1 \nu}{1 + \alpha^2 \lambda_1}$ .

**Theorem.** The pullback  $\mathcal{D}_\alpha$ -attractor  $\hat{A}_\alpha = \{A_\alpha(t) : t \in \mathbb{R}\}$  is compact in  $(H^2(\Omega))^2 \cap V$  in the sense that for any fixed  $t \in \mathbb{R}$ ,  $A_\alpha(t)$  is a compact set in  $(H^2(\Omega))^2 \cap V$ .

# Upper semicontinuity of pullback attractors for 2D Navier-Stokes-Voigt equations

- $\alpha = 0$ : Navier-Stokes equations
- Pullback attractors in  $V$  for strong solutions to 2D Navier-Stokes equations:
  - bounded domain: Real et. al., *J. Diff. Equa.* (2012).
  - unbounded domain: Anh-Trang, *preprint* (2013).
- Upper semicontinuity of pullback attractors:

**Theorem.** The family of pullback  $\mathcal{D}_\alpha$ - attractors  $\hat{\mathcal{A}}_\alpha = \{A_\alpha(t) : t \in \mathbb{R}\}$  for 2D Navier-Stokes-Voigt equations is upper semicontinuous in  $V$  at  $\alpha = 0$ , that is, for any  $t \in \mathbb{R}$ ,

$$\lim_{\alpha \rightarrow 0} \text{dist}_V(A_\alpha(t), A(t)) = 0.$$

# Existence, uniqueness and stability of stationary solutions

- Let  $f \in V'$  be independent of time. A **stationary solution** to problem (2.1) is an element  $u^* \in V$  such that

$$\nu((u^*, v)) + b(u^*, u^*, v) = \langle f, v \rangle, \quad \forall v \in V.$$

- Assume that  $f \in V'$  and

$$\nu^2 > \frac{c_0^2}{\lambda_1^{1/2}} \|f\|_*,$$

where  $c_0$  is the best constant in Ladyzhenskaya's inequality. Then there exists a unique stationary solution to problem (2.1). Moreover, this stationary solution is exponentially stable:

$$|u(t) - u^*|^2 \leq e^{-\lambda t} |u_0 - u^*|^2.$$



# Some generalizations of Navier-Stokes equations

- 3D convective Brinkman-Forchheimer equations (or tamed/damping Navier-Stokes equations)

$$\begin{cases} u_t - \nu \Delta u + (u \cdot \nabla)u + f(x, u) + \nabla p & = g, \\ \nabla \cdot u & = 0, \end{cases}$$

where  $f(\cdot, u) \sim |u|^{p-1}u, p > 3$ .

- Global attractor (bounded domain): Kalantarov-Zelik, *Comm. Pure Appl. Anal.* (2012).
- Pullback attractor (unbounded domain): Anh-Trang, *submitted*.

## Some generalizations of Navier-Stokes equations (cont.)

- Kelvin-Voigt-Brinkman-Forchheimer equations

$$\begin{cases} u_t - \nu \Delta u - \alpha^2 \Delta u_t + (u \cdot \nabla)u + f(x, u) + \nabla p & = g, \\ \nabla \cdot u & = 0, \end{cases}$$

where  $f(\cdot, u) \sim |u|^{p-1}u, p \geq 1$ .

See Anh-Trang, *Nonlinear Anal.* (2013).

- 2D second grade fluid equations

$$\begin{cases} \partial_t(u - \alpha\Delta u) - \nu\Delta u + \operatorname{curl}(u - \alpha\Delta u) \times u + \nabla p = f, & x \in \Omega, t > \tau \\ \nabla \cdot u = 0, & x \in \Omega, t > \tau, \\ u(x, t) = 0, & x \in \partial\Omega, t > \tau, \\ u(x, \tau) = u_0(x), & x \in \Omega. \end{cases}$$

Here the parameters  $\alpha$  and  $\nu$  are given positive constants and the initial datum  $u_0$  satisfies the compatibility condition

$$\operatorname{div}u_0 = 0 \text{ in } \Omega \quad \text{and} \quad u_0 = 0 \text{ on } \partial\Omega.$$

- Fluids with pressure dependent viscosities

## Some coupled systems in fluid mechanics

- the 2D Bénard problem (or Boussinesq system): a system with the Navier-Stokes equations for the velocity field coupled with a convection-diffusion equation for the temperature.

$$\begin{cases} \partial_t u + (u \cdot \nabla)u - \nu \Delta u + \nabla p = f_u(x, t) + \alpha \vec{e}_2 (T - T_r), \\ \partial_t T + (u \cdot \nabla)T - \kappa \Delta T = f_T(x, t), \\ \nabla \cdot u = 0, \end{cases}$$

See Anh-Son, *Math. Methods Appl. Sci.* (2013).

## Some coupled systems in fluid mechanics (cont.)

- 2D magnetohydrodynamic (MHD) equations: a system with the Navier-Stokes equations for the velocity field coupled with a convection-diffusion equation for the magnetic fields

$$\begin{cases} \frac{\partial u}{\partial t} + (u \cdot \nabla)u - \frac{1}{R_e} \Delta u + \nabla p + S \nabla \left( \frac{B^2}{2} \right) - S(B \cdot \nabla)B = f, \\ \frac{\partial B}{\partial t} + (u \cdot \nabla)B - (B \cdot \nabla)u + \frac{1}{R_m} \tilde{\text{curl}}(\text{curl } B) = 0, \\ \text{div } u = 0, \\ \text{div } B = 0. \end{cases}$$

- the two-phase flows: Cahn-Hilliard-Navier-Stokes system, Allen-Cahn-Navier-Stokes system, models for the nematic liquid crystal flows.

## Further readings

- 1 R. Temam, *Navier-Stokes Equations: Theory and Numerical Analysis*, 3rd edition, Amsterdam: North-Holland, 1984.
- 2 R. Temam, Some developments on Navier-Stokes equations in the second half of the 20th century, *Development of Mathematics 1950–2000*, 1049-1106, Birkhäuser, Basel, 2000.
- 3 J. Málek and K. Rajagopal, Mathematical issues concerning the Navier-Stokes equations and some of its generalizations, *Handbook of Differential Equations*. Vol. II, 371-459, Elsevier, 2005.
- 4 E. Feireisl and D. Pražák, *Asymptotic Behavior of Dynamical Systems in Fluid Mechanics*, Springer, 2010.
- 5 M. Petcu, R. Temam and M. Ziane, Some mathematical problems in geophysical fluid dynamics, *Handbook of Numerical Analysis*. Vol. XIV, 577-750, Elsevier, 2009.

**Thank you for your attention!**